

Tech report: augmented MPM for phase-change and varied materials

For our constitutive model we use $\hat{\Psi}_\mu(\mathbf{F}) = \Psi_\mu(J^{-\frac{1}{d}}\mathbf{F})$, where plasticity does not matter and is ignored for the purposes of computing these derivatives. For convenience, let $a = -\frac{1}{d}$, and μ subscripts are ignored. Then, $\hat{\Psi}(\mathbf{F}) = \Psi(J^a\mathbf{F})$. We begin by computing $\hat{\Psi}_\mu(\mathbf{F})$. We will use index notation for precision during the derivations. Differentiation by the matrix \mathbf{F}_{ij} is indicated by enclosing the index pair in parenthesis after a comma, as in $J_{,(ij)}$. Let $\mathbf{H} = \mathbf{F}^{-T}$. We begin with some preliminary derivatives for J and \mathbf{H} .

$$\begin{aligned}
 H_{ji}F_{jk} &= \delta_{ik} \\
 J_{,(ij)} &= JH_{ij} \\
 (J^a)_{,(ij)} &= aJ^{a-1}J_{,(ij)} \\
 &= aJ^{a-1}JH_{ij} \\
 &= aJ^aH_{ij} \\
 (H_{ji}F_{jk})_{,(rs)} &= 0 \\
 H_{ji,(rs)}F_{jk} + H_{ji}F_{jk,(rs)} &= 0 \\
 H_{ji,(rs)}F_{jk} &= -H_{ji}F_{jk,(rs)} \\
 H_{ji,(rs)}\delta_{jm} &= -H_{ji}\delta_{jr}\delta_{ks}H_{mk} \\
 H_{ji,(rs)} &= -H_{ri}H_{js}
 \end{aligned}$$

The derivatives of the quantity $J^a\mathbf{F}$ will occur frequently, so we begin by naming them and evaluating them.

$$\begin{aligned}
 \mathcal{B}_{kmij} &= (J^aF_{km})_{,(ij)} \\
 &= J^aF_{km,(ij)} + (J^a)_{,(ij)}F_{km} \\
 &= J^a\delta_{ik}\delta_{jm} + aJ^aF_{km}H_{ij} \\
 \mathcal{B}_{kmij}Z_{ij} &= J^a\delta_{ik}\delta_{jm}Z_{ij} + aJ^aF_{km}H_{ij}Z_{ij} \\
 \mathcal{B}_{kmij}Z_{ij} &= J^aZ_{km} + aJ^aF_{km}H_{ij}Z_{ij} \\
 \mathcal{B} : \mathbf{Z} &= J^a(\mathbf{Z} + a(\mathbf{H} : \mathbf{Z})\mathbf{F}) \\
 Z_{km}\mathcal{B}_{kmij} &= J^a\delta_{ik}\delta_{jm}Z_{km} + aJ^aF_{km}H_{ij}Z_{km} \\
 Z_{km}\mathcal{B}_{kmij} &= J^aZ_{ij} + aJ^aF_{km}Z_{km}H_{ij} \\
 \mathbf{Z} : \mathcal{B} &= J^a(\mathbf{Z} + a(\mathbf{F} : \mathbf{Z})\mathbf{H}) \\
 \mathcal{B}_{kmij,(rs)} &= (J^a(\delta_{ik}\delta_{jm} + aF_{km}H_{ij}))_{,(rs)} \\
 \mathcal{B}_{kmij,(rs)} &= (J^a)_{,(rs)}(\delta_{ik}\delta_{jm} + aF_{km}H_{ij}) + J^a(\delta_{ik}\delta_{jm} + aF_{km}H_{ij})_{,(rs)} \\
 \mathcal{B}_{kmij,(rs)} &= aJ^aH_{rs}(\delta_{ik}\delta_{jm} + aF_{km}H_{ij}) + aJ^a(F_{km,(rs)}H_{ij} + F_{km}H_{ij,(rs)}) \\
 \mathcal{B}_{kmij,(rs)} &= a\mathcal{B}_{kmij}H_{rs} + aJ^a(\delta_{kr}\delta_{ms}H_{ij} - F_{km}H_{rj}H_{is})
 \end{aligned}$$

With the operator \mathcal{B} , we can express the relationship between $\hat{\mathbf{A}} = \frac{\partial \hat{\Psi}}{\partial \mathbf{F}}(\mathbf{F})$ and $\mathbf{A} = \frac{\partial \Psi}{\partial \mathbf{F}}(J^a \mathbf{F})$.

$$\begin{aligned}\hat{\Psi}(F_{ij}) &= \Psi(J^a F_{ij}) \\ \hat{\Psi}_{,(ij)} &= \Psi_{,(km)}(J^a F_{km})_{,(ij)} \\ &= \Psi_{,(km)} \mathcal{B}_{kmij} \\ \hat{\mathbf{A}} &= \mathbf{A} : \mathcal{B}\end{aligned}$$

Finally, we relate $\mathcal{C} = \frac{\partial^2 \Psi}{\partial \mathbf{F} \partial \mathbf{F}}(J^a \mathbf{F})$ to $\hat{\mathcal{C}} = \frac{\partial^2 \hat{\Psi}}{\partial \mathbf{F} \partial \mathbf{F}}(\mathbf{F})$.

$$\begin{aligned}\hat{\Psi}_{,(ij)(rs)} &= (\Psi_{,(km)} \mathcal{B}_{kmij})_{,(rs)} \\ \hat{\Psi}_{,(ij)(rs)} &= \Psi_{,(km)(tu)} \mathcal{B}_{turs} \mathcal{B}_{kmij} + \Psi_{,(km)} \mathcal{B}_{kmij,(rs)} \\ \hat{\Psi}_{,(ij)(rs)} &= \Psi_{,(km)(tu)} \mathcal{B}_{turs} \mathcal{B}_{kmij} + a \Psi_{,(km)} \mathcal{B}_{kmij} H_{rs} + a J^a \Psi_{,(rs)} H_{ij} - a J^a \Psi_{,(km)} F_{km} H_{rj} H_{is} \\ \hat{\Psi}_{,(ij)(rs)} Z_{rs} &= \Psi_{,(km)(tu)} \mathcal{B}_{turs} \mathcal{B}_{kmij} Z_{rs} + a \Psi_{,(km)} \mathcal{B}_{kmij} H_{rs} Z_{rs} + a J^a \Psi_{,(rs)} H_{ij} Z_{rs} - a J^a \Psi_{,(km)} F_{km} H_{rj} H_{is} Z_{rs} \\ \hat{\mathcal{C}} : \mathbf{Z} &= (\mathcal{C} : (\mathcal{B} : \mathbf{Z})) : \mathcal{B} + a (\mathbf{H} : \mathbf{Z}) \mathbf{A} : \mathcal{B} + a J^a (\mathbf{A} : \mathbf{Z}) \mathbf{H} - a J^a (\mathbf{A} : \mathbf{F}) \mathbf{H} \mathbf{Z}^T \mathbf{H}\end{aligned}$$